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Length-Dependent Convergence and Saturation Behavior of Electrochemical, Linear Optical, Quadratic Nonlinear Optical, and Cubic Nonlinear Optical Properties of Dipolar Alkynylruthenium Complexes with Oligo(phenyleneethynylene) Bridges

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Abstract: The syntheses of *trans*-[Ru{4,4'-C≡CC₆H₂[2,5-(OEt)₂]C≡CC₆H₄NO₂)Cl(dppm)₂] (19), *trans*- $[Ru\{4,4',4''-C \equiv CC_6H_4C \equiv CC_6H_2[2,5-(OEt)_2]C \equiv CC_6H_4NO_2\}CI(dppm)_2]$ (20), trans- $[Ru\{4,4',4'',4''',4'''-C \equiv CC_6H_4-C \equiv CC_6+C \equiv CC_6$ $C = CC_{6}H_{2}[2,5-(OEt)_{2}]C = CC_{6}H_{2}[2,5-(OEt)_{2}]C = CC_{6}H_{4}NO_{2}CI(dppe)_{2}] (21), trans-[Ru{4,4',4'',4''',-C = CC_{6}H_{4}-C,C)]C = CC_{6}H_{4}-CC_{6}H_{4}NO_{2}CI(dppe)_{2}] (21), trans-[Ru{4,4',4'',4'',-C = CC_{6}H_{4}-C,C)]C = CC_{6}H_{4}-CC_{6}H_{4}NO_{2}CI(dppe)_{2}] (21), trans-[Ru{4,4',4'',4'',-C = CC_{6}H_{4}-C,C)]C = CC_{6}H_{4}-CC_{6}H_{4}NO_{2}CI(dppe)_{2}] (21), trans-[Ru{4,4',4'',4'',-C = CC_{6}H_{4}-C,C)]C = CC_{6}H_{4}-CC_$ $C = CC_{6}H_{2}[2,5-(OEt)_{2}]C = CC_{6}H_{2}[2,5-(OEt)_{2}]C = CC_{6}H_{4}NO_{2}Cl(dppm)_{2}] (22), trans-[Ru{4,4',4'',4'''-C} = CC_{6}H_{4}C = CC_{6}H_{$ $C \equiv CC_6H_4C \equiv CC_6H_2[2,5-(OEt)_2]C \equiv CC_6H_2[2,5-(OEt)_2]C \equiv CC_6H_4NO_2]C[(dppm)_2] (24) are reported, together and the second sec$ with those of precursor alkynes, complexes with the donor- π -bridge-acceptor formulation that affords efficient guadratic and cubic NLO compounds: the identity of **19** was confirmed by a structural study. The electrochemical properties of 19–24 and related complexes with shorter π -bridge ligands were assessed by cyclic voltammetry, and the linear optical, quadratic nonlinear optical, and cubic nonlinear optical properties were assayed by UV-vis-NIR spectroscopy, hyper-Rayleigh scattering studies at 1064 and 1300 nm, and broad spectral range femtosecond Z-scan studies, respectively. The Rull/III oxidation potentials and wavelengths of the optical absorption maxima decrease on π -bridge lengthening, until the tri(phenyleneethynylene) complex is reached, further chain lengthening leaving these parameters invariant; theoretical studies employing time-dependent density functional theory have shed light on this behavior. The quadratic nonlinearity β_{1064} and two-photon absorption cross-section reach maximal values at this same π -bridge length, a similar saturation behavior that may reflect a common importance of ruthenium-to-alkynyl ligand charge transfer in electronic and optical behavior in these molecules.

Introduction

Considerable interest has been shown in the nonlinear optical (NLO) properties of organometallic complexes,¹⁻⁶ with the majority of studies involving propagation of structure–property relationships developed in organic systems into the organometallic domain. Organic compounds with a donor– π -bridge–acceptor composition were shown to possess large quadratic NLO properties that could be enhanced by increasing

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donor or acceptor strength or by bridge modification,⁷ while compounds with extended π -systems were shown to possess

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Organometallic Complexes for Nonlinear Optics. Part 44. Part 43: Rigamonti, L.; Babgi, B.; Cifuentes, M. P.; Roberts, R. L.; Petrie, S.; Stranger, R.; Righetto, S.; Teshome, A.; Asselberghs, I.; Clays, K.; Humphrey, M. G. *Inorg. Chem.* 2009, 48, 3562.

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large cubic NLO coefficients.⁸ In proceeding to the organometallic realm, the substitution of organic donor groups by ligated metal centers increases molecular design flexibility (the metal, co-ligands, and coordination geometry can all potentially be varied); this often results in quadratic and/or cubic NLO response enhancement and in some instances permits NLO switching by oxidation/reduction at the metal.^{9–18} The most heavily studied organometallic complexes in nonlinear optics thus far have been metallocenyl^{19–62} or metal alkynyl comp-

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lexes.^{63–66} We have reported extensive studies developing structure–quadratic NLO activity relationships for dipolar metal alkynylcomplexes^{67–78} and structure–cubic NLO trends, ^{12,14,15,68,72,79–91} but almost all complexes incorporate π -bridge units containing

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one or two phenyl rings coupled together in various ways; our only study examining π -bridge lengthening, in proceeding from *trans*-[Ru(4-C=CC₆H₄NO₂)Cl(dppm)₂] to *trans*-[Ru(4,4'-C= CC₆H₄C=CC₆H₄NO₂)Cl(dppm)₂] and then *trans*-[Ru(4,4',4''-C=CC₆H₄C=CC₆H₄C=CC₆H₄NO₂)Cl(dppm)₂], resulted in a nonlinear increase in quadratic nonlinearity as measured by hyper-Rayleigh scattering (HRS) at 1064 nm [β_{1064} : 767 to 833 to 1379; β_0 : 129 to 161 to 365 (10⁻³⁰ esu)].⁶⁸ [We have very recently prepared and assessed further examples with π -bridges containing three phenyl rings linked by a variety of yne and *E*-ene groups and again noted an increase in quadratic nonlinearity.^{1,92}] Unfortunately, attempts to extend the oligo(phenyleneethynylene) series to the tetra(phenyleneethynylene)-containing example resulted in an alkynyl complex with

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minimal solubility, which hampered reaction and complex characterization and rendered solution NLO studies impossible. We report herein several new oligo(phenyleneethynylene)-containing alkynes incorporating solubilizing substituents, the corresponding ruthenium alkynyl complex derivatives, assessment of the impact of π -bridge lengthening on electrochemical and linear optical properties, theoretical studies using time-dependent density functional theory (TD-DFT) directed at rationalizing our experimental observations, quadratic nonlinear optical properties from hyper-Rayleigh scattering studies at two wavelengths, and wide-spectral-range wavelength-dependence studies of the cubic nonlinear optical properties of selected complexes from Z-scan studies; these studies demonstrate that electrochemical and optical properties in this system converge at the tri(phenyleneethynylene) π -bridge length.

Experimental Section

Materials. All reactions were performed under a nitrogen atmosphere with the use of Schlenk techniques unless otherwise stated. Tetrahydrofuran (THF) was dried by distilling over sodium/ benzophenone; all other solvents were used as received. Petrol is a fraction of boiling range 60-80 °C. Chromatography was performed on silica gel or ungraded basic alumina. Ethynyltrimethylsilane (Wacker), tert-butyllithium (concentration determined prior to reaction by titration against diphenylacetic acid), 1,2diiodoethane, 4-bromo-1-iodobenzene, bromine, iodine, sodium hexafluorophosphate, tetra-n-butylammonium fluoride, copper(I) iodide, sodium sulfite, magnesium sulfate, sodium acetate, sodium bicarbonate, potassium iodate, PdCl₂(PPh₃)₂, and Pd(PPh₃)₄ (Aldrich) were used as received. The following were prepared by literature procedures: 1,4-diethoxybenzene,⁹³ 1,4-dibromo-2,5diethoxybenzene,⁹⁴ 4-Me₃SiC=CC₆H₄I, 4-HC=CC₆H₄NO₂,⁹⁵ 4-HC \equiv CC₆H₄C \equiv CSiPrⁱ₃,⁹⁶ 4,4'-IC₆H₄C \equiv CC₆H₄C \equiv CSiPrⁱ₃,⁹⁷ cis-[RuCl₂(dppe)₂], cis-[RuCl₂(dppm)₂],⁹⁸ trans-[Ru(4-C \equiv CC₆H₄-NO₂)Cl(dppe)₂] (25),⁹⁹ trans-[Ru(4,4'-C=CC₆H₄C=CC₆H₄NO₂)- $Cl(dppe)_2$] (26), trans-[Ru(4,4',4''-C=CC_6H_4C=CC_6H_4C=CC_6H_4-C=CC_6H_4C=CC_6H_4-C=CC_6H_4C=CC_6H_4-C=CC_6H_4C=CC_6H_4-C=CC_6H_4C=CC_6H_4-C=CC_6H_4C=CC_6H_4C=CC_6H_4-C=CC_6H_4C=CC_6H_4-C=CC_6H_4C=CC_6H_4-C=CC_6H_4C=CC_6H_4-C=CC_6H $NO_2)Cl(dppe)_2]$ (27),¹ *trans*-[Ru(4-C=CC₆H₄NO₂)Cl(dppm)₂] (28),¹⁰⁰ trans-[Ru(4,4'-C=CC₆H₄C=CC₆H₄NO₂)Cl(dppm)₂] (29), *trans*-[Ru(4,4',4"-C=CC₆H₄C=CC₆H₄C=CC₆H₄NO₂)Cl(dppm)₂] (30).⁶⁸ The syntheses of 1–18 are given in the Supporting Information.

Methods and Instrumentation. Microanalyses were carried out at the Australian National University. UV—vis spectra of solutions in 1 cm quartz cells were recorded as dichloromethane solutions using a Cary 5 spectrophotometer; bands are reported in the form frequency (cm⁻¹) [extinction coefficient ($10^4 \text{ M}^{-1} \text{ cm}^{-1}$)]. Infrared spectra were recorded as solutions in dichloromethane or as KBr disks using a Perkin-Elmer System 2000 FT-IR; data are reported in cm⁻¹. ¹H (300 MHz), ¹³C (75 MHz), and ³¹P NMR (121 MHz) spectra were recorded using a Varian Gemini-300 FT NMR spectrometer and are referenced to residual chloroform (7.26 ppm), CDCl₃ (77.0 ppm), or external H₃PO₄ (0.0 ppm), respectively; atom

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labeling follows the numbering scheme in Chart S1. Electrospray ionization mass spectra (ESI-MS) were recorded using a Micromass/ Water's LC-ZMD single quadrupole liquid chromatograph MS, high-resolution (HR) ESI mass spectra were carried out utilizing a Bruker Apex 4.7T FTICR-MS instrument, and electron impact mass spectra (EI MS) were recorded using a VG Quattro II triple quadrupole MS; all mass spectrometry peaks are reported as m/z(assignment, relative intensity). Cyclic voltammetry measurements were recorded using a MacLab 400 interface and MacLab potentiostat from ADInstruments. The supporting electrolyte was 0.1 M (NBuⁿ₄)PF₆ in distilled, deoxygenated CH₂Cl₂. Solutions containing ca. 1×10^{-3} M complex were maintained under argon. Measurements were carried out at room temperature using Pt disk working, Pt wire auxiliary, and Ag/AgCl reference electrodes, such that the ferrocene/ferrocenium redox couple was located at 0.56 V (peak separation ca. 0.09 V). Scan rates were typically 100 mV $\ensuremath{s^{-1}}$.

Synthesis of trans-[Ru{4,4'-C=CC₆H₂[2,5-(OEt)₂]C=CC₆H₄-NO₂}Cl(dppm)₂] (19). cis-[RuCl₂(dppm)₂] (132 mg, 0.14 mmol) and NaPF₆ (31.3 mg, 0.17 mmol) were added to a solution of 5 (47.1 mg, 0.14 mmol) in CH₂Cl₂ (30 mL). The yellow mixture was stirred at room temperature overnight. NEt₃ (1 mL) was added and the red mixture stirred at room temperature for 6 h. The reaction mixture was purified by passing through a short pad of alumina, eluting with CH_2Cl_2 /petrol/NEt₃ (10:10:1). Reduction in volume of the solvent on a rotary evaporator afforded 19 as a dark red powder (120 mg, 69%). ESI-MS: 1240 $([M]^+, 100)$, 1205 $([M - Cl]^+, 10)$. Anal. Calcd for C70H60ClNO4P4Ru • 0.25CH2Cl2: C, 66.92; H, 4.84; N, 1.11. Found: C, 66.95; H, 5.11; N, 1.25. UV-vis: 19 600 [1.8], 26 500 [2.2], 30 400 [1.9]. IR (CH₂Cl₂): 2200, 2064 (C≡C). ¹H NMR: δ 1.25, 1.41 (2 t, 6H, $J_{\rm HH}$ = 7 Hz, 2CH₃), 3.67, 3.82 (2 q, 4H, J_{HH} = 7 Hz, 2CH₂), 4.80, 5.32 (2 m, 4H, 2PCH₂), 5.19, 6.77 (2 s, 2H, H₁₆, H₁₉), 5.28 (s, 0.5H, CH₂Cl₂), 6.99-7.60 (m, 44H, H_{24} , Ph), 8.20 (d, $J_{HH} = 8$ Hz, 2H, H_{25}). ¹³C NMR: δ 15.1, 15.2 (CH₃), 47.6 (PCH₂), 63.7, 64.4 (CH₂), 91.6, 94.2 (C₂₁, C₂₂), 104.2, 110.4 (C₁₅, C₁₈), 114.1, 117.4 (C₁₆, C₁₉), 123.1 (C₂₃), 123.6 (C₂₅), 127.5 (Ph), 129.0 (d, J_{CP} = 36 Hz, Ph), 131.6 (C₂₄), 133.5 (d, J_{CP} = 38 Hz, Ph), 146.2 (C_{26}), 152.4, 153.7 (C_{17} , C_{20}), C_{13} , C_{14} not observed. ³¹P NMR: δ -4.7. The syntheses of **20**-**24** are similar and are given in the Supporting Information.

X-ray Structural Study of trans-[Ru{4,4'-C=CC₆H₂[2,5-(OEt)₂]- $C = CC_6H_4NO_2$ [Cl(dppm)₂] (19). A crystal suitable for the structural study was grown by slow diffusion of methanol into a dichloromethane solution of 19 at -20 °C. A unique diffractometer data set was collected on a Nonius Kappa CCD diffractometer using the ω scan technique (graphite-monochromated Mo K α radiation, 0.71073 Å). The unit cell parameters were obtained by least-squares refinement¹⁰¹ of 13 727 reflections with $2.6^{\circ} \le \theta \le 27.5^{\circ}$ at 200 K. An analytical absorption correction was applied, using numerical methods,¹⁰² implemented from MAXUS,¹⁰³ and equivalent reflections were merged. The structure was solved by direct methods, expanded using Fourier techniques,¹⁰⁴ and refined using the CRYSTALS software package.^{105,106} One of the phenyl rings attached to a dppm moiety was found to be disordered, and so each atom was split. The total site occupancies were set to be one for each set of split atoms. All other non-hydrogen atoms were refined anisotropically. Hydrogen atoms were included in the refinement

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at idealized positions, riding on the atoms to which they are bonded, and their positions were frequently recalculated. Conventional residuals R and R_w on |F| were 4.3% and 15.2%, respectively; a Chebyshev weighting function was employed. The largest peaks in the final difference electron density map are located near the Ru atom. Molecular graphics were displayed using the PLATON software.107

Crystal/Refinement Data: $C_{70}H_{60}CINO_4P_4Ru$, M = 1239.67, monoclinic, $P2_1/c$, a = 22.2245(2) Å, b = 12.7995(1) Å, c =23.5634(2) Å, $\beta = 116.4624(5)^\circ$, V = 6000.62(9) Å³. D_c (Z = 4) = 1.327 g cm⁻³. μ_{Mo} = 0.46 mm⁻¹; specimen: 0.40 × 0.25 × 0.24 mm; $T_{min/max}$ = 0.857/0.916. $2\theta_{max}$ = 55°; N_{total} = 125 417 (CCD diffractometer, monochromatic Mo K α radiation, $\lambda = 0.71073$ Å; T 200 K) merging to 13 727 unique ($R_{int} = 0.044$), $N_o = 3517$ (I $\geq 3\sigma(I)$ refining to $R = 0.0425, R_w = 0.152.$

Theoretical Studies. Calculations were performed using the Amsterdam Density Functional (ADF) package ADF2006.01,¹⁰⁸ developed by Baerends et al.^{109,110} These calculations were undertaken to characterize the lowest frequency allowed singlephoton transitions of a set of model compounds containing the oligo(phenyleneethynylene) bridges of compounds 19 to 24. These models were of the form $[Ru](C_2C_6H_4)_iC_2C_6H_4NO_2$, where [Ru] =trans-RuCl(PH₂CH₂PH₂)₂ and i = 1-4. In the discussion that follows, the models are denoted as 19M (i = 1), 20M (i = 2), **23M** (i = 3), and **24M** (i = 4) to indicate the laboratory compounds of which they are structural homologues. Note that, for reasons of computational expediency, the solubilizing ethoxy substituents of the laboratory compounds have been omitted in our model calculations. The presence or absence of alkoxy substituents on the phenylene rings has been found to exert only a minor influence on the intensity and wavelength of major single-photon transitions in a detailed computational study of related compounds.⁹⁷ Omission of the ethoxy substituents in this instance allowed the structural symmetry of the models to be constrained to C_{2v} throughout. In all calculations and for all atoms, the Slater-type orbital basis sets used were of triple- ζ -plus-polarization quality (TZP). Electrons in orbitals up to and including 1s {C, N, O}, 2p {P, Cl}, and 4d {Ru} were treated in accordance with the frozen-core approximation. Geometry optimizations employed the gradient algorithm of Versluis and Ziegler.¹¹¹ Functionals used in the optimization calculations were the local density approximation (LDA) to the exchange potential,¹¹² the correlation potential of Vosko, Wilk, and Nusair (VWN),¹¹³ and the nonlocal corrections of Perdew, Burke, and Enzerhof (PBE).¹¹⁴ Following optimization of the model compounds, timedependent density functional theory (TD-DFT) calculations were pursued using either PBE or the asymptotically correct functional of van Leeuwen and Baerends (LB94).

HRS Measurements. An injection-seeded Nd:YAG laser (Qswitched Nd:YAG Quanta Ray GCR, 1064 nm, 8 ns pulses, 10 Hz) was focused into a cylindrical cell (7 mL) containing the sample. The intensity of the incident beam was varied by rotation of a half-wave plate placed between crossed polarizers. Part of the laser pulse was sampled by a photodiode to measure the vertically polarized incident light intensity. The frequency-doubled light was collected by an efficient condenser system and detected by a

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photomultiplier. The harmonic scattering and linear scattering were distinguished by appropriate filters; gated integrators were used to obtain intensities of the incident and harmonic scattered light. The absence of a luminescence contribution to the harmonic signal was confirmed by using interference filters at different wavelengths near 532 nm. All measurements were performed in THF using *p*-nitroaniline ($\beta = 21.4 \times 10^{-30}$ esu¹¹⁵) as a reference. Solutions were sufficiently dilute that absorption of scattered light was negligible.

For studies at 1.300 μ m, a Tsunami-pumped OPAL (model Spectra-Physics) was used. With a high repetition rate of the laser, high-frequency demodulation of fluorescence contributions can be effected, a full description being given in ref 116. All measurements were performed in THF using Disperse Red 1 (DR1, $\beta = 54 \times 10^{-30}$ esu in chloroform) as a reference. Experiments utilized low chromophore concentrations, the linearity of the HRS signal as a function of the chromophore concentration confirming that no significant self-absorption of the SHG signal occurred.

Z-Scan Studies. The compounds were dissolved in dichloromethane at concentrations in the range 0.5-1.2% w/w, and the solutions were placed in 1 mm path length Starna glass cells, stoppered, and sealed with Teflon tape. Due to experimental limitations, only one concentration of each compound was examined. The Z-scans were carried out using a femtosecond laser system composed of a Clark-MXR CPA-2001 regenerative amplifier acting as a 775 nm pump and a Light Conversion TOPAS optical parametric amplifier. The system was operated at a repetition rate of 250 Hz to minimize the influence of thermal effects and photochemical decomposition. The experiments were carried out at a number of wavelengths in the range 530-1500 nm. To obtain the relevant wavelengths, the optical parametric amplifier was tuned appropriately and one of the following modes of output was selected: idler-pump mixing, signal doubling, idler doubling, or the signal. The unwanted components of the TOPAS output were discarded through the use of polarizing optics, color glass filters, and spatial filtering.

The laser beam was attenuated to energies in the μ J/pulse range and directed through a standard Z-scan setup equipped with a beam splitter, allowing one to record open-aperture and closed-aperture Z-scans simultaneously. The beam was focused so as to provide a spot size in the range $w_0 \approx 40-80 \,\mu\text{m}$, ensuring that the Rayleigh range $z_{\rm R} = \pi w_0^2 / \lambda$ was always larger than the total thickness of the sample (\approx 3 mm, which includes two glass walls and the solution inside the cell). The cell travel was z = -40 to 40 mm, and the data were recorded using three Si or InGaAs photodiodes monitoring the input pulse energy, the open-aperture signal, and the closedaperture signal, respectively. The outputs were fed into three channels of a boxcar averager, which was GPIB-interfaced with a data collection computer. The shapes of the closed-aperture scans, open-aperture scans, and the curves obtained by dividing a closedaperture curve by an open-aperture curve were analyzed with the help of a custom fitting program that used equations derived by Sheik-Bahae et al.¹¹⁷ to calculate the theoretical curves. In this way, the values of the real and imaginary part of the nonlinear phase shift $\Delta \Phi_0$ (corresponding to the refractive and absorptive nonlinearity, respectively) could be obtained. At each wavelength the energy of the laser pulses was adjusted, based on closed-aperture Z-scans obtained for 1 mm cells filled with the solvent alone and/ or scans for a 3 mm thick silica glass plate. Typically, we required $\Delta \Phi_0$ for such scans to be about 0.5–1 rad. The nonlinear refractive index of silica, n_2 , is well known across a wide wavelength range, so the real part of the nonlinear phase shift can be used to calculate the light intensity. From the comparison of n_2 values for silica compiled by Milam¹¹⁸ we employed

$$n_{2,\text{silica}}(\bar{\nu}) = 2.82 \times 10^{-16} - 3 \times 10^{-21} \bar{\nu} + 2 \times 10^{-25} \bar{\nu}^2 (\text{cm}^2/\text{W})$$

for the spectral dependence of n_2 for silica, where $\bar{\nu} = 1/\lambda$ is the wavenumber expressed in cm⁻¹. This equation is a good approximation for values of $n_{2,\text{silica}}$ in the range of $\lambda = 500-1500$ nm. Note that most of our previous papers assumed $n_{2,\text{silica}} = 3 \times 10^{-16} \text{ cm}^2/\text{W}$ independent of the wavelength, which is also a reasonably good approximation.

We employed intensities on the order of 100 GW/cm². Since our measurements were always uniformly calibrated to the nonlinearity of silica glass, this avoided the need for detailed information of beam geometry and pulse shape.

The real and imaginary parts of the second hyperpolarizability, γ , of the solutes were computed assuming additivity of the nonlinear contributions of the solvent and the solute and the applicability of the Lorentz local field approximation.¹¹⁹ The values of the two-photon absorption (TPA) cross-section σ_2 were computed from the absorptive part of the nonlinearity determined from the Z-scans. In all cases, errors of the relevant quantities were estimated from the assessed accuracies of the parameters for the fitting of Z-scans for the solutions and the corresponding scans for the solvent.

Results

Synthesis and Characterization of Alkynes. The acetylenes required for the alkynyl complex syntheses were prepared by extensions of established organic synthetic procedures (Schemes 1–6). Monobromination of 1,4-di(ethoxy)benzene using bromine and sodium acetate in acetic acid afforded 1-bromo-2,5diethoxybenzene (1), in a procedure previously applied to afford the 1-bromo-2,5-di(*n*-propoxy)benzene analogue (Scheme 1).¹²⁰ Subsequent iodination with iodine/potassium iodate in an acetic acid/carbon tetrachloride mixture afforded 4-bromo-2,5-di ethoxy-1-iodobenzene (2) in good yield. This facile synthesis of 2 is preferable to the trans-halogenation of 1,4-dibromo-2,5di(ethoxy)benzene employing *tert*-butyllithium and 1,2-diiodoethane.

Both the bromo and iodo functionalities in **2** are reactive. 4-Ethynyl-1-nitrobenzene and 4-ethynyl-1-(triisopropylsilyl)ethynylbenzene reacted with **2** under Sonogashira conditions at room temperature to give 2,5-(EtO)₂-4-BrC₆H₂C \equiv C-4-C₆H₄NO₂ (**3**) and 4-Pr^{*i*}₃SiC \equiv CC₆H₄C \equiv CC₆H₂-2,5-(OEt)₂-4-Br (**6**), respectively, in excellent yields (Scheme 2). The former reacted with ethynyltrimethylsilane under Sonogashira conditions to afford 4-Me₃SiC \equiv C-2,5-(EtO)₂C₆H₂-1-C \equiv CC₆H₄-4-NO₂ (**4**) in fair yield. The diyne **4** was desilylated to afford 4-HC \equiv C-2,5-(EtO)₂C₆H₂-1-C \equiv CC₆H₄-4-NO₂ (**5**).

Syntheses of tri(phenyleneethynylene)-containing compounds are summarized in Scheme 3. **5** underwent sequential Sonogashira coupling with 4-bromo-1-iodobenzene and then ethynyltrimethylsilane to afford $4-O_2NC_6H_4C\equiv CC_6H_2-2,5-(OEt)_2-4-C\equiv CC_6H_4-4-C\equiv CSiMe_3$ (**7**), which was desilylated by tetra*n*-butylammonium fluoride, giving $4-O_2NC_6H_4C\equiv CC_6H_2-2,5-(OEt)_2-4-C\equiv CC_6H_4-4-C\equiv CH(8)$. A similar sequential Sonogashira coupling with 4-Br-2,5-(EtO)_2C_6H_2I and then ethynyltrimethylsilane gave $4-O_2NC_6H_4C\equiv CC_6H_2-2,5-(OEt)_2-4-C\equiv CSiMe_3$ (**10**), which was similarly desilylated to

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Scheme 1. Syntheses of 1 and 2



Scheme 2. Syntheses of 3-6



Scheme 3. Syntheses of 7-11



 $4-O_2NC_6H_4C \equiv CC_6H_2-2,5-(OEt)_2-4-C \equiv CC_6H_2-2,5-(OEt)_2-4-C \equiv CH$ (11).

Syntheses of tetra(phenyleneethynylene)-containing compounds in which two rings bear solubilizing groups are summarized in Scheme 4. $4-Pr_3SiC \equiv CC_6H_4C \equiv CC_6H_2-2,5-(OEt)_2-4-C \equiv CC_6H_2-2,5-(OEt)_2-4-C \equiv CC_6H_4-4-NO_2$ (12) is available from reaction of 5 with 6 under Sonogashira conditions; 12 was desilylated by tetra-*n*-butylammonium fluoride to give 4-HC≡CC₆H₄C≡CC₆H₂-2,5-(OEt)₂-4-C≡CC₆H₂-2,5-(OEt)₂-4-C≡CC₆H₄-4-NO₂ (**14**) in quantitative yield. **14** is also accessible from a sequential Sonogashira coupling of **11** with 4-bromo-1-iodobenzene and then ethynyltrimethylsilane, affording 4-Me₃-SiC≡CC₆H₄C≡CC₆H₂-2,5-(OEt)₂-4-C≡CC₆H₂-2,5-(OEt)₂-4-C≡CC₆H₄-4-NO₂ (**13**) and a subsequent desilylation. Syntheses of tetra(phenyleneethynylene)-containing compounds in which one ring bears solubilizing groups are summarized in Scheme

Scheme 4. Syntheses of 12-14



Scheme 5. Syntheses of 15-17



Scheme 6. Synthesis of 18



5. Sequential Sonogashira coupling of **8** with 4-bromo-1iodobenzene and then ethynyltrimethylsilane gave 4-Me₃SiC \equiv CC₆H₄-1-C \equiv CC₆H₄-4-C \equiv CC₆H₂-2,5-(OEt)₂-4-C \equiv CC₆H₄-4-NO₂ (**15**) in a good yield, subsequent desilylation affording 4-HC \equiv CC₆H₄C \equiv CC₆H₄-4-C \equiv CC₆H₂-2,5-(OEt)₂-4-C \equiv CC₆H₄-4-NO₂ (**17**). Alternatively, **5** was coupled to 4-IC₆H₄C \equiv CC₆H₄-4-C \equiv CC₆H₂-2,5-(OEt)₂-4-C \equiv CC₆H₄-4-NO₂ (**16**), which was desilylated to give **17** in a better yield. The penta(phenyleneethynylene) 4-Prⁱ₃SiC \equiv CC₆H₄-1-C \equiv CC₆H₄-4-C \equiv CC₆H₂-2,5-(OEt)₂-4-C \equiv CC₆H₄-4-C \equiv CC₆H₄-4-C \equiv CC₆H₄-2,5-(OEt)₂-4-C \equiv CC₆H₄-4-NO₂ (**18**) was obtained via Sonogashira coupling of **11** with 4-IC₆H₄C \equiv CC₆H₄-4-C \equiv CSiPrⁱ₃ in a fair yield (Scheme 6).

Synthesis and Characterization of Alkynyl Complexes. The new alkynyl complexes **19–24** were prepared using established methodologies (Schemes 7 and 8)^{77,99,121–131} and characterized



by a combination of IR, ¹H, ¹³C, and ³¹P NMR spectroscopies and ESI mass spectrometry. The IR spectra contain characteristic ν (C=C) bands at 2063–2074 cm⁻¹ for the metal-bound alkynyl group, with the higher values corresponding to complexes with the longer π -bridges (22–24). The ³¹P NMR spectrum of the dppe-containing complex 21 contains one singlet resonance at 49.8 ppm, confirming the *trans*-disposed diphosphine ligands. The spectra of the dppm-containing complexes similarly contain a singlet; the ³¹P resonance for the complex with ethoxy substituents on the ring adjacent to the metal center (19: –4.7 ppm) is somewhat downfield of those of other dppm-containing complexes (–5.5 ppm). Mass spectra contain molecular ions and/or fragment ions corresponding to loss of a chloro ligand.

The identity of **19** was confirmed by a single-crystal X-ray diffraction study. Figure 1 contains an ORTEP plot showing the molecular geometry, together with selected bond lengths

Scheme 7. Syntheses of 19-23



 $[Ru] = trans-RuCl(dppe)_2 (21)$ [Ru] = trans-RuCl(dppm)_2 (19, 20, 22, 23)

Scheme 8. Synthesis of 24



and angles. Bond lengths and angles are similar to those of the structurally characterized analogues trans-[Ru(C=C-4-C₆H₄C=CC₆H₄-4-X)Cl(dppm)₂] (X = H, ⁶⁸ N=NC₆H₄-4-NO₂¹²⁴), the only notable exception being Cl1-Ru1-Cl 170.33(9)° (**19**), cf. 175.9(2)° (X = H), 175.1(17)° (X = N=NC₆H₄-4-NO₂), a distortion in **19** likely to derive from the need to minimize unfavorable steric interactions between the C4-ethoxy group and the P-phenyl rings. The dihedral angle between the C3/C8 and C11/C16 phenyl planes in **19** (5.7(2)°) is much smaller than those in precedent structural studies (26.75° (X = H), 21.9° (X = N=NC₆H₄-4-NO₂), 33.9° (*trans*-[Ru(C=C-4-C₆H₄C=C-4-C₆H₄C=CPh)₂(dmpe)₂]¹³²)). Because

coplanar phenyl rings in extended π -systems enhance π -delocalization, this structural difference is likely to favorably influence nonlinearity in solid-state materials; however, the centrosymmetric packing in the crystal structure of **19** will result in no bulk second-order susceptibility.

Electrochemical and Linear Optical Studies. Electrochemical properties of ruthenium alkynyl complexes are of significant interest (see ref 92 and references therein). The results of cyclic voltammetric studies of the new ruthenium alkynyl complexes **19–24** are collected in Table 1, together with data from related complexes **25–30**.



Figure 1. Molecular geometry and atomic labeling scheme for *trans*-[Ru{4,4'-C=CC₆H₂[2,5-(OEt)₂]C=CC₆H₄NO₂}Cl(dppm)₂] (**19**). Selected bond lengths (Å) and angles (deg): Ru1–P1 2.3291(8), Ru1–P2 2.3362(8), Ru1–P3 2.3613(8), Ru1–P4 2.3504(8), Ru1–Cl1 2.5176(7), Ru1–Cl 1.988(3), C1–C2 1.217(5), C2–C3 1.421(4); P1–Ru1–P2 71.76(3), P3–Ru1–P4 70.87(3), C11–Ru1–Cl 170.33(9), Ru1–C1–C2 177.2(3), C1–C2–C3 175.8(4).

Table 1. Cyclic Voltammetric Data for Complexes 19–30 ^a				
complex	$E^{0}_{ m ox}$ (V) $[i_{ m po}/i_{ m pa}$	Ru ^{II/III} E ⁰ red	(V) [<i>i</i> _{po} / <i>i</i> _{pa}] NO ₂ ^{0/-1}	ref
trans-[Ru(4-C \equiv CC ₆ H ₄ NO ₂)Cl(dppe) ₂] (25)	0.74 [0.	- [6	-0.84 [0.8]	68
<i>trans</i> -[Ru(4,4'-C=CC ₆ H ₄ C=CC ₆ H ₄ NO ₂)Cl(dppe) ₂] (26)	0.60 [1]		-0.94 [0.9]	1
trans-[Ru(4,4',4''-C=CC ₆ H ₄ C=CC ₆ H ₄ C=CC ₆ H ₄ NO ₂)Cl(dppe) ₂] (27)	0.58 [1]		-0.91[0.9]	1
<i>trans</i> -[Ru {4,4',4'',4'', C=CC ₆ H ₄ C=CC ₆ H ₂ [2,5-(OE) ₂]C=CC ₆ H ₂ [2,5-(OE) ₂]C=CC ₆ H ₄ NO ₂ } C1(dppe) ₂] (21)	0.56 [0.	- [6	-1.16 [0.9]	this work
<i>trans</i> -[Ru(4-C \equiv CC ₆ H ₄ NO ₂)Cl(dppm) ₂] (28)	0.72 [1]		-0.81 [0.7],	68
<i>trans</i> -[Ru(4,4'-C=CC ₆ H ₄ C=CC ₆ H ₄ NO ₂)Cl(dppm) ₂] (29)	0.57 [0.	- [6	-0.90 [0.7]	68
<i>trans</i> -[Ru {4,4'-C=CC ₆ H ₂ [2,5-(OEb ₂]C=CC ₆ H ₄ NO ₂ }Cl(dppm) ₂] (19)	0.50 [1]		-0.91[0.9]	this work
trans-[Ru(4,4',4''-C=CC ₆ H ₄ C=CC ₆ H ₄ C=CC ₆ H ₄ NO ₂)Cl(dppm) ₂] (30)	0.54 [1]		-0.86 [0.9]	68
$trans-[Ru \{4,4,''-C=CC,H_4C=CC,H_2[2,5-(OEt),]C=CC,H_4NO_3]C[C,H_4NO_3]C[C,H_4NO_3]C[(dppm),] (20)$	0.54^{1}		-0.86 [0.9]	this work
$nrans-[Ru \{4,4',4'',4''',-C=CC_6H_4C=CC_6H_3[2,5-(OEt),]C=CC_6H_3[2,5-(OEt),]C=CC_6H_4NO_3[C(dppm)_3] $ (22)	0.55 [1]		-0.87 [0.9]	this work
<i>trans</i> -[Ru[4,4',4'',-C=CC,H_C=CC,H_AC=CC,H_AC=CC,H_AC=CC,H_ANO, ACI(dpom),1 (23)	0.55 [1]		-0.89 [0.9]	this work
<i>trans</i> -[Ru {4,4',4'',4''',4''',4'''',C=CC ₆ H ₄ C=CC ₆ H ₄ C=CC ₆ H ₂ [2,5-(OEt) ₂]C=CC ₆ H ₂ [2,5-(OEt) ₂]C=CC ₆ H ₄ NO ₂ }CI(dppm) ₂] (24)	0.56 [0.	- [6	-0.90 [0.6]	this work
Table 2. Experimental Linear Optical, Quadratic Nonlinear Optical, ^a and Two-Photon Absorption ^b Response Parameter Image: Non-Photon Absorption ^b Response Parameter Table 2. Experimental Linear Optical, Quadratic Nonlinear Optical, ^a and Two-Photon Absorption ^b Response Parameter Image: Non-Photon Absorption ^b Response Parameter Table 2. Experimental Linear Optical, Quadratic Nonlinear Optical, ^a and Two-Photon Absorption ^b Response Parameter Image: Non-Photon Absorption ^b Response Parameter Table 2. Experimental Linear Optical, Quadratic Nonlinear Optical, ^a and Two-Photon Absorption ^b Response Parameter Image: Non-Photon Absorption ^b Response Parameter Table 2. Experimental Linear Optical, Quadratic Nonlinear Optical, ^a and Two-Photon Absorption ^b Response Parameter Image: Non-Photon Absorption ^b Response Parameter Trans-Ru(4-C=CC,HLOS)(dippe)] (25) Image: Ru(4,4''-C=CC,HLOS)(dippe)] (26) Image: Ru(4,4''-C=CC,HLOS)(dippe)] (27) Trans-Ru(4,4''-C=CC,HLOS)(dippe)] (28) Image: Ru(4,4''-C=CC,HLOS)(dippe)] (28) Image: Ru(4,4'',4'''-C=CC,HLOS)(dippe)] (28) Trans-Ru(4,4''-C=CC,HLOS)(dippe)] (28) Image: Ru(4,4'',4'',-C=CC,HLOS)(dippe)] (28) Image: Ru(4,4'',4'',-C=CC,HLOS)(dippe)] (28) Trans-Ru(4,4'',4'',-C=CC,HLOS)(dippe)] (28) Image: Ru(4,4'',4'',-C=CC,HLOS)(dippe)] (28) Image: Ru(4,4'',4'',4'',4'',4'',4'',4'',4'',4'',4	$ \begin{array}{c} \mbox{m} [f_{\rm s} & \beta_{\rm rusa} \left(10^{-30} \\ \mbox{M}^{-1} \right) \\ \mbox{M}^{-1} & \beta_{\rm rusa} \left(10^{-30} \\ \mbox{m}^{\rm sol} \right) \\ \mbox{m}^{-1} & \mbox{sub} \\ \mbox{m}^{\rm sol} & \mbox{m}^{\rm sol} \\ \mbox{m}^{\rm sol} & \$	$ \begin{array}{c} \beta_0 \\ (10^{-30} \\ (10^$	$ \begin{array}{c} \begin{array}{c} \mbox{$P_{\rm C}$} & \mbox{$P_{\rm C}$} \\ \beta_0 & \mbox{$h_{\rm max}$} \\ (10^{-30} & \mbox$} \\ (10^{-30} & \mbox{$h_{\rm max}$} \\$	ref 1 1 1 1 1 1 1 3, 1, this work 68, 1, this work this work this work this work this work
^a Conditions: measurements were carried out in THF; all complexes are optically transparent at 1064 and 1300 nm. The β dat out in CH ₂ Cl ₂ . ^c Corrected for resonance enhancement at 532 nm using the two-level model with $\beta_0 = \beta[1 - (2\lambda_{max}/1064)^2][1 - (\lambda_{max}/1300)^2]$. ^c The accuracy of the σ_2 data are $\pm 10\%$.	are $\pm 10\%$ unless s - $(\lambda_{\text{max}}/1064)^2$]. ^d	pecified otherwise. ^b Corrected for resona	Conditions: measurer ince enhancement at	nents were carried 650 nm using the



Figure 2. Optical spectra of *trans*- $[Ru(4,4'-C \equiv CC_6H_4C \equiv CC_6H_4NO_2)-Cl(dppm)_2]$ (29) and *trans*- $[Ru\{4,4'-C \equiv CC_6H_2[2,5-(OEt)_2]C \equiv CC_6H_4-NO_2\}Cl(dppm)_2]$ (19).



Figure 3. Optical spectra of *trans*-[Ru(4,4',4''-C \equiv CC₆H₄C \equiv CC₆H₄C \equiv CC₆H₄NO₂)Cl(dppm)₂] (**30**) and *trans*-[Ru{4,4',4''-C \equiv CC₆H₄C \equiv CC₆H₂[2,5-(OEt)₂]C \equiv CC₆H₄NO₂}Cl(dppm)₂] (**20**).



Figure 4. Optical spectra of *trans*-[Ru{4,4',4'',4'''-C \equiv CC₆H₄C \equiv CC₆H₂[2,5-(OEt)₂]C \equiv CC₆H₂[2,5-(OEt)₂]C \equiv CC₆H₄NO₂}Cl(dppm)₂] (**22**) and *trans*-[Ru { 4, 4', 4'', 4''' - C \equiv C C ₆ H ₄ C \equiv C C ₆ H ₄ C \equiv C C ₆ H ₂ [2, 5 - (OEt)₂]C \equiv CC₆H₄NO₂}Cl(dppm)₂] (**23**).

As with other *trans*-bis(bidentate diphosphine)ruthenium monoalkynyl complexes, the cyclic voltammograms of the new complexes contain a reversible or quasi-reversible anodic wave assigned to the Ru^{II/III} oxidation process, in the range 0.50-0.56 V, and a less reversible cathodic wave assigned to the NO₂^{0/-1} reduction process, in the range -0.86 to -1.16 V. These data, together with those of the previously reported complexes **25**–**30**, permit comment on the effect of π -bridge lengthening and solubilizing group incorporation upon metal-localized oxidation potential. For the dppe-containing complexes, OPE increase (proceeding from **25** to **26** to **27** and finally **21**) leads to a progressive decrease in Ru^{II/III} oxidation potential, but this is not clear-cut because of the introduction of ethoxy groups at **21**. The effect of these solubilizing substituents can in principle



be deconvoluted by examination of the more extensive data available for the dppm-containing complexes. Thus, introduction of solubilizing substituents at the π -bridge phenyl ring adjacent to the metal (proceeding from 29 to 19) results in a significant increase in the ease of oxidation (ΔE_{ox}^0 0.07 V), whereas solubilizing group introduction more remote from the metal (proceeding from 30 to 20) and further increase in solubilizing groups (proceeding from 23 to 22) results in no change in oxidation potential; these results, which suggest electron density at the metal is particularly sensitive to adjacent ring modification, are consistent with those seen upon nitro group incorporation into such complexes (across the series 19-30, oxidation potential is significantly higher for complexes in which the nitro group is attached to the ring adjacent to the metal, namely, 25 and 28). These results suggest that the effect of π -bridge lengthening on ease of oxidation should only be assessed with complexes incorporating the same number of solubilizing groups. With this in mind, the present data suggest π -bridge lengthening (proceeding from 28 to 29 and then 30, from 20 to 23, and from 22 to 24) progressively increases the ease of oxidation until the tri(phenyleneethynylene) is reached, after which the oxidation potential is invariant.

Absorption maxima and intensities from electronic spectra are collected in Table 2, and the evolution in optical spectra on systematic structural modification is illustrated in Figures 2-5. Incorporating solubilizing groups on the ring adjacent to the metal (proceeding from 29 to 19) results in a red-shift and decrease in intensity of the absorption bands (Figure 2), while introducing the solubilizing groups at a more remote ring (proceeding from 30 to 20) results in a decrease in intensity of the higher energy bands (Figure 3); a further increase in number of solubilizing groups remote from the metal (proceeding from 23 to 22) results in a slight increase in intensity of bands, but again little shift in the location of the absorption maxima (Figure 4). Minimizing solubilizing substituents while π -bridge lengthening, in proceeding from 28, through 29, 30, and 23, to 24 (Figure 5), results in a distinct blue-shift from the mono(phenyleneethynylene) 28 to the tri(phenyleneethynylene) 30, with little further change in v_{max} , but an increase in absorptivity, in proceeding to the penta(phenyleneethynylene) 24. Closer inspection of the spectra reveals that the next-higher-energy bands red-shift and gain in intensity on chain-lengthening, to eventually coalesce with the lowest energy band for 23 and 24.

Table 3. Computed Significantly Allowed Single-Photon Transitions (those with oscillator strengths exceeding f = 0.3 au) for [Ru](C₂C₆H₄),C₂C₆H₄NO₂ (i = 1-4), Obtained through PBE/TZP and LB94/TZP//PBE/TZP Calculations

		PBE/TZP							LB94/TZP					
model complex ^a	symm ^b	nc	E/eV ^d	f/au ^e	OCC. ^f	virt. ^f	wt % ^g	nc	<i>E</i> /eV ^d	f/au ^e	occ. ^f	virt. ^f	wt % ^g	
19M	A_1	1	1.533	0.43	${}^{14}B_2$	$^{15}B_{2}$	94	1	1.167	0.35	${}^{14}B_2$	¹⁵ B ₂	92	
	A_1	3	2.810	0.81	${}^{14}B_2$	$^{16}B_2$	74	3	2.441	0.67	${}^{14}B_2$	${}^{16}B_2$	44	
20M	A_1	1	1.112	0.33	${}^{17}B_2$	${}^{18}B_2$	97	1	0.749	0.30	${}^{17}B_2$	${}^{18}B_2$	96	
	A_1	3	2.162	0.83	${}^{17}B_2$	${}^{19}B_2$	88	3	1.870	0.78	${}^{17}B_2$	${}^{19}B_2$	65	
	A_1	5	2.975	1.01	${}^{14}B_2$	${}^{18}B_2$	42	5	2.527	0.34	${}^{14}B_2$	${}^{18}B_2$	65	
	A_1							7	2.801	0.70	${}^{17}B_2$	${}^{20}B_2$	70	
	A_1							10	3.206	0.36	${}^{15}B_2$	${}^{19}B_2$	49	
23M	A_1	3	1.755	0.69	${}^{20}B_2$	${}^{22}B_2$	93	3	1.507	0.62	${}^{20}B_2$	${}^{22}B_2$	84	
	A_1	6	2.568	0.46	${}^{17}B_2$	${}^{21}B_2$	47	5	2.086	0.34	${}^{17}B_2$	${}^{21}B_2$	75	
	A_1	5	2.495	0.84	${}^{20}B_2$	${}^{23}B_2$	75	6	2.244	0.20	${}^{20}B_2$	${}^{23}B_2$	79	
	A_1	7	2.657	0.48	${}^{19}B_2$	${}^{22}B_2$	56	7	2.326	0.84	${}^{19}B_2$	${}^{22}B_2$	73	
	A_1	8	2.947	0.52	${}^{18}B_2$	${}^{22}B_2$	57	9	2.695	0.76	${}^{18}B_2$	${}^{22}B_2$	53	
24M	A_1	3	1.514	0.32	${}^{22}B_2$	${}^{24}B_2$	95	2	1.034	0.13	${}^{22}B_2$	${}^{24}B_2$	96	
	A_1	2	1.490	0.49	${}^{23}B_2$	${}^{25}B_2$	95	3	1.264	0.41	${}^{23}B_2$	${}^{25}B_2$	94	
	A_1	4	1.942	0.38	${}^{21}B_2$	${}^{24}B_2$	82	4	1.393	0.10	${}^{21}B_2$	${}^{24}B_2$	85	
	A_1	6	2.282	0.74	${}^{20}B_2$	${}^{24}B_2$	64	5	1.764	0.28	${}^{20}B_2$	${}^{24}B_2$	79	
	A_1	5	2.112	0.48	${}^{23}B_2$	${}^{26}B_2$	84	6	1.876	0.38	${}^{23}B_2$	${}^{26}B_2$	89	
	A_1	7	2.346	0.76	${}^{22}B_2$	${}^{25}B_2$	66	7	2.064	0.88	${}^{22}B_2$	${}^{25}B_2$	80	
	A_1	8	2.620	0.98	${}^{21}B_2$	${}^{25}B_2$	44	9	2.374	0.90	${}^{21}B_2$	${}^{25}B_2$	42	
	A_1	10	2.797	0.15	${}^{22}B_2$	$^{26}B_2$	19	12	2.620	1.01	${}^{22}B_2$	${}^{26}B_2$	56	

^{*a*} Notation used for calculation on model compounds is consistent with that indicated in the main text. ^{*b*} Symmetry classification of the identified transition within the indicated symmetry classification. ^{*d*} Calculated transition energy in electronvolts, at the indicated level of theory. ^{*f*} Calculated transition oscillator strength, in atomic units, at the indicated level of theory. ^{*f*} Principal pair of occupied and virtual orbitals involved in the identified transition. ^{*g*} Percentage contribution of principal-orbital character to the calculated transition, at the indicated level of theory.

Theoretical Studies. Our TD-DFT calculations on the four model compounds [Ru]($C_2C_6H_4$) $_iC_2C_6H_4$ NO₂ (i = 1-4) delivered, for each model, the 50 lowest energy symmetry-allowed single-photon transitions. Of this number, only a handful of calculated transitions for each model (all such transitions being of A₁ symmetry for the $C_{2\nu}$ -constrained models) have expectation values *f* greater than about 0.3 atomic unit (see Table 3).

Orbital characteristics for models **19M**, **20M**, **23M**, and **24M** show considerable common features, as is to be expected from their structural similarity. The HOMO in each case $(^{11+3i}B_2)$ is dominated by acetylenic character concentrated on the C₂ unit immediately adjacent to Ru, while the LUMO $(^{12+3i}B_2)$ is characterized by C=N bonding of the terminal phenylene to the nitro group. The next few lower lying B₂-symmetry occupied orbitals show acetylenic character at progressively greater separation from Ru at increasing distance from the "frontier", while the next few higher lying B₂-symmetry virtual orbitals

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are cumulenic in character, associated with C_2 units that are closer to the Ru as the orbital energy increases above the LUMO.

Perhaps because the spatial separation of domains for the HOMO and LUMO becomes more exaggerated as chain length *i* is increased, the LUMO—HOMO excitation ${}^{1}A_{1}$ becomes progressively less important with increasing *i*, with the calculated *f* value for this transition decreasing monotonically at both the PBE/TZP and LB94/TZP levels of theory. Instead, there is a marked tendency toward the strongest transitions being those where there is a close spatial correspondence between the occupied and virtual domains.

Analysis of the TD-DFT results is hampered by inconsistencies between the intensities calculated at the PBE/TZP level of theory and those obtained at the LB94/TZP level.¹³³ For example, for 24M (see Figure 6) the most intense transitions at the PBE/TZP level of theory are ${}^{6}A_{1}$ (${}^{24}B_{2} \leftarrow {}^{20}B_{2}$), f = 0.74 au; ${}^{7}A_{1}$ (${}^{22}B_{2}$ $- {}^{25}B_{2}$), f = 0.76 au; and ${}^{8}A_{1}$ (${}^{21}B_{2}$ $- {}^{25}B_{2}$), f = 0.98au. For the same compound according to LB94/TZP, the latter two transitions (here ${}^{7}A_{1}$ and ${}^{9}A_{1}$) are again calculated to be intense, with f values of 0.88 and 0.90, very close to the PBE/ TZP results. However, the transition involving ${}^{24}B_2 \leftarrow {}^{20}B_2$ character, here ${}^{5}A_{1}$, has a much lower intensity (f = 0.28), while the most intense transition according to LB94/TZP, A1 $(^{22}B_2 \leftarrow ^{26}B_2), f = 1.01$ au, has only a weak counterpart (f = 0.15) in the PBE/TZP calculations. While the two TD-DFT methods employed yield somewhat inconsistent results for the identities of the most intense single-photon transitions, it remains notable that the most intense transitions consistently display

⁽¹³³⁾ The problems of selecting the most appropriate functional for predicting absorption spectra using TD-DFT have been discussed recently. See, for example: (a) Jacquemin, D.; Perpète, E. A.; Scuseria, G. E.; Ciofini, I.; Adamo, C. J. Chem. Theory Comput. 2008, 4, 123. (b) Peach, M. J. G.; Benfield, P.; Helgaker, T.; Tozer, D. J. J. Chem. Phys. 2008, 128, 044118. (c) Peach, M. J. G.; Le Sueur, C. R.; Ruud, K.; Guilaume, M.; Tozer, D. J. Phys. Chem. Chem. Phys. 2009, 11, 4465.



Figure 6. Energy-level diagram (obtained at the PBE/TZP level of theory) and orbital diagram for calculated intense single-photon transitions of the model compound **24M**. Transition term values are shown for the PBE/TZP and LB94/TZP levels of theory, respectively; a term value in parentheses indicates that the transition in question is not notably intense (f < 0.3 au) at the indicated level of theory. A common feature of these intense transitions is the generally close spatial correspondence between the electronic domains of the relevant occupied (solid line) and virtual (dotted line) orbitals. All of the transitions identified as intense have excitation energies, at the PBE/TZP level of theory, of 2.0 eV (\sim 16 000 cm⁻¹) or higher. In contrast, the HOMO–LUMO gap is much less, only 0.73 eV (\sim 6000 cm⁻¹), but does not feature significant overlap between occupied and virtual domains and is calculated to contribute only weakly ($f \sim 0.1$ au) to the single-photon excitation spectrum of **24M**.

strong domain-region overlap between the identified occupied and virtual orbitals. An apparent consequence of the increasing importance of this effect with increasing chain length is that the frequencies of the most intense transitions do not decrease as rapidly as does the HOMO/LUMO energy separation. Indeed, as the crucial orbitals are displaced progressively more from the frontier, it may well appear that the important transitions become blue-shifted on chain lengthening.

A second mechanism also exists whereby the most prominent transitions detected can appear to blue-shift on chain lengthening. While the lowest energy conformer for **24M**, as for all of its smaller congeners, features coplanarity of all of the coaxial aryl rings, the barriers to intramolecular rotation along any of the ethynyl spindles are not large, with the calculated rotational transition states lying typically 10 kJ mol⁻¹, or less, above the all-coplanar conformer.¹³⁴ Further, while the primary effect of moderate on-axis aryl rotation (i.e., for dihedral angle values of up to 30°) is generally to reduce the energy of each allowed transition by a modest quantity (i.e., by 300 cm⁻¹ or less), there is a more prominent effect on the calculated oscillator strength for each transition, with internal rotation dampening some

transitions and strengthening others. This phenomenon can lead to apparent blue-shifting on chain lengthening (since increasing chain length corresponds to an increase in the number of opportunities for on-axis aryl rotation) when rotation enhances the relative strength of a higher energy transition relative to its lower energy neighbor. Since the spectral features observed in the laboratory are typically broad (~5000 cm⁻¹ fwhm or similar), a rotationally mediated change in the relative intensity of two or more close-lying transitions may well manifest as the blue-shifting of a single broad peak effectively enveloping several transitions.

It is important to note that, while both of these phenomena can result in blue-shifting on chain lengthening, neither effect is exactly systematic: they therefore appear to be of little predictive value. However, they provide a rationalization for the somewhat counterintuitive observation that elongation of the phenyleneethynylene "axis" can result in an increase in the frequency of the observed spectral peak.

Quadratic Nonlinear Optical Studies. The quadratic nonlinearities of **19–24** have been determined at 1064 nm; these results, together with previously reported measurements for the



Figure 7. Comparison of the dispersion of the complex hyperpolarizability of (a) trans-[Ru(4-C=CC₆H₄NO₂)Cl(dppm)₂] (28), (b) trans-[Ru(4,4'-C=CC₆H₄C=CC₆H₄NO₂)Cl(dppm)₂] (29), (c) trans-[Ru(4,4',4''-C=CC₆H₄C=CC₆H₄C=CC₆H₄NO₂)Cl(dppm)₂] (20), (d) trans-[Ru(4,4',4'',4'''-C=CC₆H₄C=CCC₆H₄C=C

related complexes 25–30, are presented in Table 2. The major focus of the present study is to assess the impact of π -bridge lengthening upon optical properties. For the dppe-containing complexes, π -bridge lengthening, in proceeding from 25 to 26, 27, and finally 21, results in nonlinearity peaking at the tri(phenyleneethynylene)-containing complex 27, but there are insufficient data to deconvolute the influence of the ethoxy substituents introduced into 21 to ensure sufficient solubility. We have a more extensive set of dppm-containing complexes that permit further comment. Introduction of solubilizing groups (proceeding from 29 to 19, or 30 to 20) results in a ca. 30% reduction in β_{1064} , while further increasing the ethoxy content (proceeding from 23 to 22) results in no further change, within the error margins. There is sufficient data to offer a comment on the effect on β of π -bridge lengthening; proceeding from 28 to 29, and then 30 (no ethoxy substituents), 20 to 23 (two ethoxy substituents), and 22 to 24 (four ethoxy substituents) suggests that β peaks at the tri(phenyleneethynylene) stage, after which a saturation effect is apparent.

Table 2 also contains a more limited set of data at 1300 nm and two-level-corrected values for the two wavelengths. The results for the former are in contrast to those at 1064 nm. Introducing solubilizing substituents, in proceeding from **29** to **19**, results in an increase in nonlinearity, while a further increase



Figure 8. Comparison of the two-photon absorption dispersion of (a) trans-[Ru(4-C=CC₆H₄NO₂)Cl(dppm)₂] (28), (b) trans-[Ru(4,4'-C=CC₆H₄C=CC₆H₄C=CC₆H₄O₂)Cl(dppm)₂] (29), (c) trans-[Ru(4,4',4'',-C=CC₆H₄C=CC₆H₄C=CC₆H₄O₂)Cl(dppm)₂] (20), (d) trans-[Ru(4,4',4'',4'''-C=CC₆H₄C=CC₆H₄C=CC₆H₄C=CC₆H₄O₂)Cl(dppm)₂] (20), (d) trans-[Ru(4,4',4'',4'''-C=CC₆H₄C=CC₆C₆C=CC₆H₄C=CC₆C+CC₆C=CC₆H₄C=CC₆C+CC₆C=CCC₆C=CC₆C=CC₆C=CCC₆C

in ethoxy content (proceeding from 23 to 22) results in a decrease in β value. The π -bridge lengthening does not result in a saturation point being reached, β increasing on proceeding from 28 to 29 (no ethoxy substituents) and 20 to 23 (two ethoxy groups). Two-level-corrected data should be viewed with extreme caution; while the trends are broadly similar to those for the corresponding β_{λ} values, the shortcomings of the model have been discussed extensively elsewhere (ref 100 and refs therein).

Cubic Nonlinear Optical Studies. The cubic NLO properties of selected complexes were assessed over a wide spectral range (Figure 7), the specific complexes (**20**, **23**, **24**, **28**, **29**) being chosen to ensure sufficient solubility for the Z-scan studies. Figure 7 contains the hyperpolarizability values plotted against twice the measurement wavenumber, to facilitate comparison with the linear optical spectrum included in the figure. The studies were undertaken in essentially transparent regions of the UV-vis-NIR spectra (from 6700 cm⁻¹ to the onset of one-

photon absorption). A few data were collected with moderate linear absorption, and these consequently have large error factors.

With reference to the results depicted in Figure 7, there is generally an approximate correspondence between the one-photon absorption bands and the main two-photon absorption band. The low-energy nonlinear absorption band shown in Figure 7 in the range ca. $12000-18000 \text{ cm}^{-1}$, and appearing for laser wavelengths of ca. 1100-1500 nm, seems to have a multiphoton character with possible strong contributions from three-photon absorption.

Figure 8 shows the absorptive part of the cubic nonlinearity of the five compounds replotted as the (effective) two-photon cross-section versus wavelength of the incident laser beam, with relevant data being collected in Table 2. The rapid increase of the maximal value of the two-photon cross-section in proceeding from **28**, through **29**, to **20** (at ~900 nm for all the compounds)

is followed by stabilization of the value in proceeding to the longer π -bridge compounds 23 and 24.

Discussion

The present studies have afforded a series of OPEs endfunctionalized by ligated Ru and nitro groups. Increasing OPE length is associated with decreasing solubility, which at the tetra(phenyleneethynylene) length is sufficiently diminished as to cause problems with chemical synthesis. Alkoxy substituents have been employed to ensure adequate solubility both for chemical synthesis and for NLO studies. Solubilizing groups such as these are not necessarily electronically and optically innocent, and so, recognizing that these groups are essential, we have sought to define their influence. Introducing solubilizing groups at the phenylene adjacent to the metal results in a decrease in E_{ox}^0 and a red-shift in the low-energy absorption bands. We have not probed this modification theoretically, but introduction of electron-donating substituents at the bridge phenyl ring adjacent to the ruthenium is expected to increase electron richness at the metal center, and thereby increase ease of oxidation. This substitution will also destabilize the HOMO and SHOMO, which are localized at the metal and adjacent ring, resulting in a red-shift in the optical absorption maximum. Introducing solubilizing groups at rings remote from the metal center results in a much smaller red-shift in the optical absorption maximum and does not affect the ease of oxidation, consistent with the aforementioned HOMO/SHOMO localization. Incorporation of these solubilizing groups results in a ca. 30% reduction in β_{1064} and β_0 . The calculations are consistent with the dominant charge-transfer transition (which would be expected to contribute significantly to the observed nonlinearity) being metal-to-bridge in character and with a diminished importance of this transition on solubilizing group incorporation and thereby a decreased nonlinearity.

The π -bridge length-dependent evolution of optical properties in organic oligomers has been of considerable interest.^{134,135} Poly(phenyleneethynylene)s have potential as photoluminescence or nonlinear optical materials or as fluorescent chemosensors, and monodisperse examples of a wide range of lengths have been prepared.^{136–138} Related to the present work, a series of OPEs end-functionalized by di(dodecyl)amino and nitro groups has been examined, with the linear optical absorption maximum blue-shifting on chain-length increase until the trimer is reached, further π -system lengthening leaving λ_{max} invariant.¹³⁹ Third-harmonic generation studies of these dialkylaminonitro-substituted OPEs were hampered by solubility problems.¹⁴⁰ The present studies expand the scope of donor-acceptorfunctionalized OPEs to include examples incorporating a ligated metal center. The blue-shift of λ_{max} on chain-lengthening this metal-containing system mirrors that seen in the examples above. Our theoretical studies have identified two possible sources of the observed blue-shift of λ_{max} on chain-lengthening: decreasing importance of the LUMO \leftarrow HOMO transition and on-axis aryl rotation-mediated change in the relative intensities of two or more close-lying transitions.

The major outcome of the present work has been to define the effective saturation length of OPE examples end-functionalized by ligated metal and nitro group (donor-bridge-acceptor arylalkynylruthenium complexes), for which metal-centered oxidation potential, linear optical absorption maximum, quadratic optical nonlinearity assessed by HRS at 1064 nm, and two-photon-absorption cross-section assessed by Z-scan all plateau at the tri(phenyleneethynylene) complex. These studies have therefore afforded the specific bridge length for the most efficient optical materials with this metal-containing donor group, as well as providing an interesting comparison of saturation length with related, purely organic, compounds.

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Supporting Information Available: List of authors of ref 108, the syntheses of compounds 1-18 and complexes 20-24, and a cif file giving solution and refinement details and tables of atomic coordinates, bond lengths, and bond angles for the X-ray structural study of 19. This material is available free of charge via the Internet at http://pubs.acs.org.

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